

# An Eight-Device Extended-Resonance Power-Combining Amplifier

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**Abstract**—An extended resonance technique for designing power amplifiers is investigated. This technique provides a planar power-combining structure whose primary distinction from conventional structures is that it does not require quarter-wave or multiple quarter-wave spacing between devices. A theoretical analysis of extended-resonance combiners is presented, followed by experimental data from an eight-device power amplifier. This approach is adaptable to monolithic circuit-fabrication techniques and should be useful for millimeter-wave applications.

## I. INTRODUCTION

THE greatest need for power combining is at microwave and millimeter-wave frequencies where an individual solid-state device is often incapable of producing adequate power for most radar and communication systems. To remedy this, power combiners are used to obtain sufficient power from solid-state devices. Typical power combiners employ Wilkinson or hybrid structures in which devices are interconnected with quarter-wave transmission lines to achieve power combining [1]–[7]. These structures become quite large when many devices are employed. Such large combining structures often have losses that rival the device gains, limiting their use at high frequencies.

An extended-resonance technique was introduced and applied to power-combining oscillators in [8]. A four-device power-combining amplifier based on extended resonance was later presented in [9]. Extended resonance effectively places devices in shunt to combine the power from each device. The magnitude of the voltage at each device is the same, and the gain of the extended-resonance amplifier is equal to the gain of a single-device amplifier. The significance of extended resonance lies in its ability to either divide power between, or combine power from, devices without the need for quarter-wave or half-wave transmission lines. Shorter line lengths may be used. This enables a compact circuit to be designed that is particularly suitable for monolithic microwave integrated circuit (MMIC) implementation. This method may also be used in conjunction with conventional methods to enhance their performance.

In this paper, a comprehensive study of extended-resonance power-combining amplifiers is discussed. This study develops the theoretical power-combining mechanism of extended

resonance along with experimental results from a new eight-device power amplifier at X-band. This eight-device amplifier handles twice as much power as the four-device amplifier presented in [9]. In addition, the eight-device amplifier presented here utilizes shorter transmission lines and maintains the same bandwidth achieved with the four-device amplifier.

The extended-resonance amplifier is superficially similar in structure to a distributed amplifier [3], but is actually quite different in its performance. The distributed amplifier is a traveling-wave structure, while the extended-resonance amplifier is actually a power-combining circuit; each device amplifying the same magnitude excitation. In addition, the gain of the extended-resonance multiple-device amplifier is the same as that for a single-device amplifier (assuming lossless transmission lines) regardless of the number of devices used in the circuit. This is indeed not the case for distributed amplifiers.

The authors have recently become aware of previous work on “ladder-type” microwave power-divider/combiner structures in rectangular waveguides [10], [11] that is somewhat similar to work reported herein. In [10] and [11], separate three-port passive power-divider and power-combiner systems in rectangular waveguide were implemented. However, when designing power amplifiers based on this technique, certain constraints are placed on the design of the dividing and combining circuits. A design technique that addresses these constraints is presented in this paper. Furthermore, an eight-device extended-resonance amplifier is constructed using microstrip.

## II. EXTENDED-RESONANCE POWER-COMBINING AMPLIFIERS

Fig. 1 illustrates one version of an extended-resonance power-combining amplifier utilizing n-FET's. Both the gates and drains are sequentially linked with transmission lines. The electrical lengths of the transmission lines connecting each device are denoted by  $\theta g_m$  and  $\theta d_m$  for the gate lines and the drain lines, respectively, where  $1 \leq m \leq n - 1$ . The input and output are located at the gate of the first device and the drain of the nth device, respectively. It is assumed that each device has the same gate admittance  $Yg = Gg + jBg$  and drain admittance  $Yd = Gd + jBd$ . The dividing and combining circuits may be considered separately, as shown in Figs. 2 and 3. Here, the gate and drain admittances may be assumed to be the simultaneous conjugate match admittances. In Fig. 2, line length  $\theta g_{n-1}$  transforms  $Y = Gg + jBg$  from the nth device, to its conjugate  $Y^* = Gg - jBg$  at the location of the next device. Adding the gate admittance of

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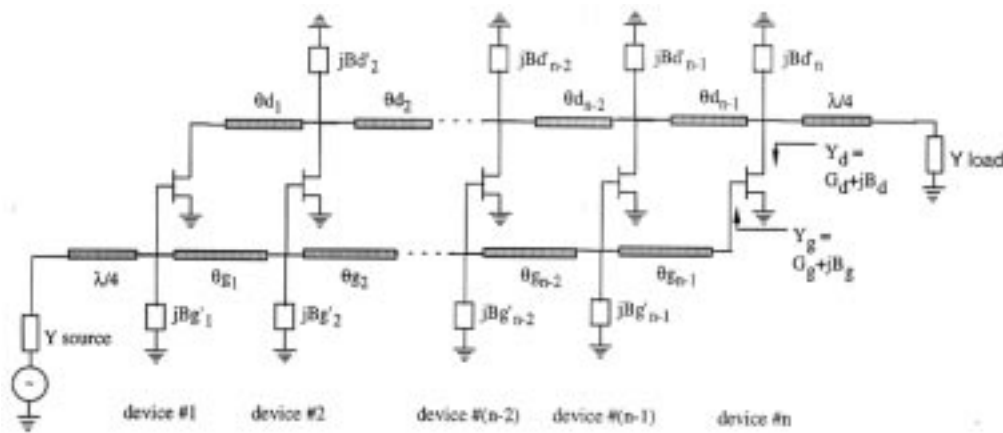
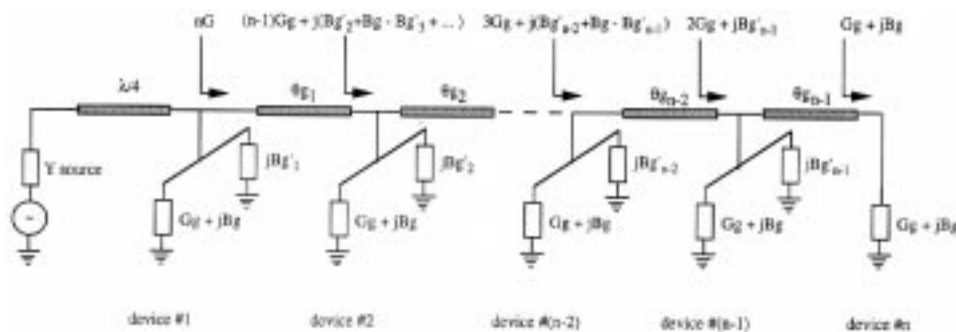
Fig. 1. An  $n$ -device extended-resonance amplifier.

Fig. 2. Input circuit.

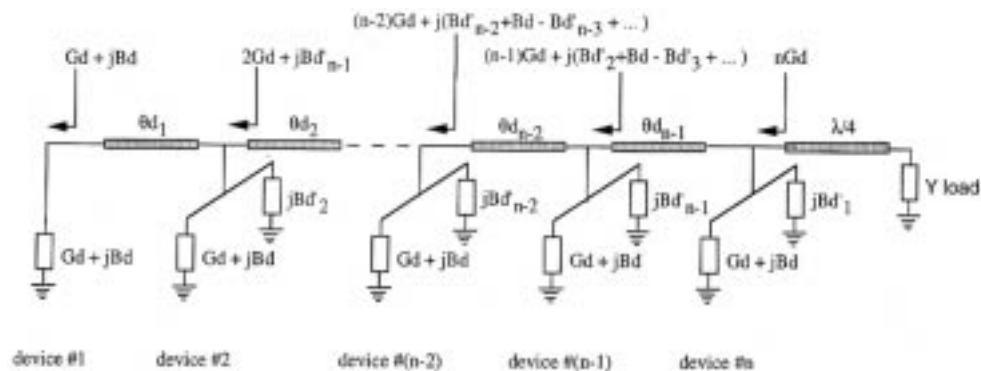


Fig. 3. Output circuit.

the next device, the imaginary components cancel; hence, the gate admittance at device  $n - 1$  is  $2G_g$ . A shunt susceptive element,  $jBg'_{n-1}$  may be placed at this device so that line length  $\theta_{g_{n-2}}$  transforms  $Y = 2G_g + jBg'_{n-1}$  from the second device to its conjugate  $Y^* = 2G_g - jBg'_{n-1}$  at the next device (the choice of  $jBg'_{n-1}$  is, at this point, completely arbitrary). This process continues all the way to the gate of the first device. The shunt susceptance,  $jBg'_1$  can be chosen to cancel the total reactive component at the input, and the input admittance  $nG_g$  can be matched to a given source admittance using a quarter-wave transformer. Similarly, in the drain circuit of Fig. 3, length  $\theta_{d_1}$  transforms  $Y = G_d + jB_d$  from the first device to  $Y^* = G_d - jB_d$  at the second device. Line length

$\theta_{d_2}$  transforms  $Y = 2G_d + jBd'_{n-1}$  at the second device to  $Y = 2G_d - jBd'_2$  at the third device. This continues to the drain of the  $n$ th device, where the leftover susceptance is canceled by  $jBd'_n$ .

The basic design procedure for an extended-resonance power amplifier will now be considered. It will be shown that an input signal applied to the gate of the first FET will be divided equally among all FET's. Each FET amplifies  $1/n$  of the input power and delivers it to the output combining circuit where the power from each FET is recombined at the load. Assuming lossless lines, it can also be shown that the total output power is equal to  $n$  times the power generated by each device.

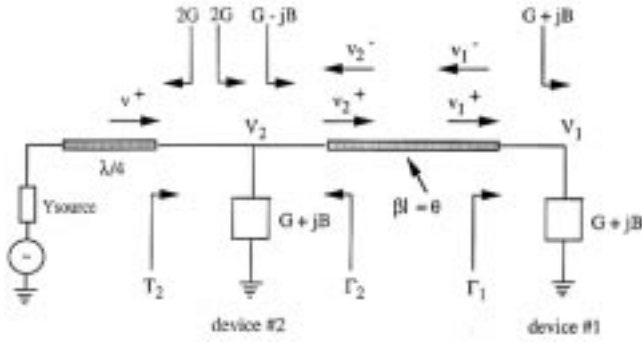


Fig. 4. Analysis of input circuit.

### III. ANALYSIS OF THE GATE EXTENDED-RESONANCE CIRCUIT

For simplicity, consider the two-device extended-resonance circuit shown in Fig. 4. Two devices, each represented by a normalized admittance  $G + jB$ , are connected with a transmission line of electrical length  $\theta$  and characteristic admittance  $Y_o$ . An analysis of this circuit can be done by determining the forward and reverse steady-state traveling waves along the transmission line after a source is applied to the input. Referring to Fig. 4,  $\Gamma_1$  is the reflection coefficient looking to the right at the first device,  $\Gamma_2$  is the reflection coefficient looking to the left at the second device, and  $T_2$  is the transmission coefficient looking to the right at the second device.  $\Gamma_1$ ,  $\Gamma_2$ , and  $T_2$  can be expressed in terms of  $Y_o$ ,  $G$ , and  $B$  as

$$\Gamma_1 = \frac{Y_o - G - jB}{Y_o + G + jB} \quad (1)$$

$$\Gamma_2 = \frac{Y_o - 3G - jB}{Y_o + 3G + jB} \quad (2)$$

$$T_2 = \frac{2Y_o}{Y_o + 2G} \quad (3)$$

In the following equations, it will be convenient to express  $\Gamma_1$  in terms of its magnitude and phase angle  $\gamma$ :

$$\Gamma_1 = |\Gamma_1|e^{j\gamma}$$

The voltage  $V_1$  across the terminals of the first device is expressed in terms of the forward and reverse traveling waves at port 1:

$$V_1 = v_1^- + v_1^+ \quad (4)$$

Similarly, for  $V_2$ ,

$$V_2 = v_2^- + v_2^+ \quad (5)$$

$V_1$  and  $V_2$  can be rewritten as

$$V_1 = \frac{T_2[1 + \Gamma_1]e^{-j\theta}}{1 - \Gamma_1\Gamma_2e^{-j2\theta}}v^+ \quad (6)$$

$$V_2 = \frac{T_2[1 + \Gamma_1e^{-j2\theta}]}{1 - \Gamma_1\Gamma_2e^{-j2\theta}}v^+ \quad (7)$$

where  $v^+$  is the incident wave at the input (see Fig. 4). From (6) and (7), the ratio of  $V_1$  and  $V_2$  is

$$\frac{V_1}{V_2} = \frac{[1 + \Gamma_1]e^{-j\theta}}{1 + \Gamma_1e^{-j2\theta}} \quad (8)$$

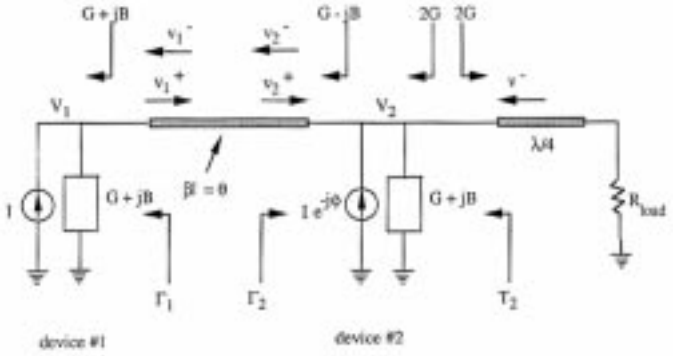


Fig. 5. Analysis of output circuit.

The length of (lossless) transmission line  $\theta$  required to transform  $Y = G + jB$  to its conjugate  $Y^* = G - jB$  can be expressed in terms of the angle  $\gamma$  of the reflection coefficient associated with the admittance  $G + jB$  as

$$\begin{aligned} \theta &= \gamma, & \text{when } B < 0 \\ \theta &= \gamma + \pi, & \text{when } B > 0. \end{aligned}$$

When  $B < 0$ , (8) becomes

$$\frac{V_1}{V_2} = \frac{[1 + |\Gamma_1|e^{j\gamma}]e^{-j\gamma}}{1 + |\Gamma_1|e^{-j\gamma}} \quad (9)$$

The magnitude of the voltage at each device is equal ( $|V_1|/|V_2| = 1$ ) and  $V_1$  lags  $V_2$  by  $\phi$  where

$$\phi = \gamma - 2 \tan^{-1} \left[ \frac{|\Gamma_1| \sin \gamma}{1 + |\Gamma_1| \cos \gamma} \right] \quad (10)$$

When  $B > 0$ , (8) becomes

$$\frac{V_1}{V_2} = \frac{[1 + |\Gamma_1|e^{j\gamma}]e^{-j(\gamma+\pi)}}{1 + |\Gamma_1|e^{-j\gamma}} \quad (11)$$

Here again, the magnitude of the voltage at each device is the same, and now  $V_1$  lags  $V_2$  by

$$\phi = \pi + \gamma - 2 \tan^{-1} \left[ \frac{|\Gamma_1| \sin \gamma}{1 + |\Gamma_1| \cos \gamma} \right] \quad (12)$$

Equations (10) and (12) are useful for determining the phase shift between devices in a two-device extended-resonance power-dividing circuit. Determining the phase angle  $\phi$  only requires knowledge of the transmission line length  $\theta$  and the reflection coefficient ( $\Gamma_1$ ) at the location of the first device. For an extended-resonance circuit containing  $n$  devices, the phase shift  $\phi$  between any two devices can be calculated by simply knowing the line lengths and the total reflection coefficient at each device port. For example, to determine the phase shift  $\phi$  between device  $m$  and device  $m+1$  in an extended-resonance circuit having an arbitrary number of devices, simply substitute  $\Gamma_m$  and  $\theta_m$  into the above equations for  $\Gamma_1$  and  $\theta$ , respectively, where  $\Gamma_m$  is the total reflection coefficient at device  $m$ , and  $\theta_m$  is the electrical length of the transmission line connecting device  $m$  to device  $m+1$ .

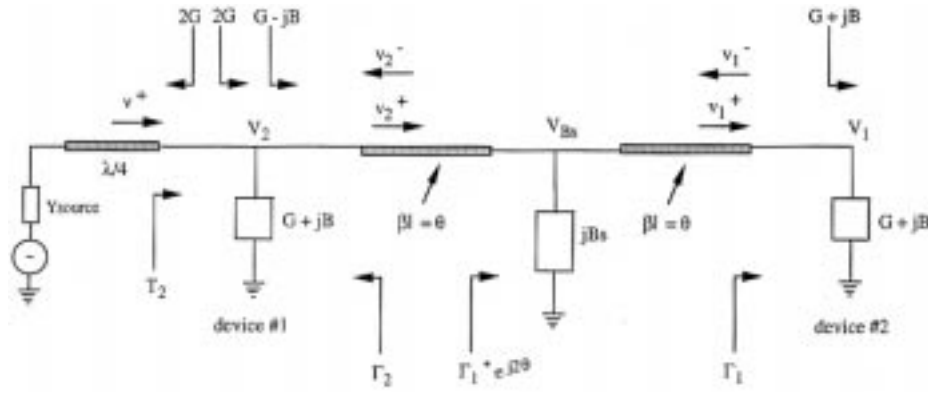


Fig. 6. Analysis of the modified input circuit.

#### IV. ANALYSIS OF THE DRAIN EXTENDED-RESONANCE CIRCUIT

Consider now the power combiner circuit in Fig. 5, which represents the drain circuit of the power amplifier. The drain of each FET is modeled as a current source and its output admittance. Based on the power-dividing analysis for the gate circuit, the current sources have equal magnitudes, with a phase difference  $\phi$  imposed by the gate circuit divider. Once again, (1) and (2) are used to express  $V_1$  and  $V_2$ , as shown in (13) and (14), at the bottom of the page. The electrical length  $\theta$  is once again chosen to be the length of line required to transform  $G + jB$  to  $G - jB$ . It can, therefore, be expressed in terms of the reflection angle  $\gamma$  as was done in the power-dividing scheme. After some mathematical manipulations, the ratio of  $V_1$  to  $V_2$  can be expressed as

$$\frac{V_1}{V_2} = \cos \phi + j \frac{G}{Y_o} \sin \theta. \quad (15)$$

Requiring the condition

$$\frac{G}{Y_o} \sin \theta = \sin \phi \quad (16)$$

to be satisfied by (15) results in  $V_1$  and  $V_2$  having equal magnitudes with a phase difference  $\phi$ . Satisfying (16) results in equal voltage at each device terminal and, consequently, total power being sent to the load.

#### V. DESIGN CONSIDERATIONS

An extended-resonance amplifier is designed using (10), (12), and (16). The value of  $\phi$ , which is calculated from the gate circuit equation (10) or (12), must be inserted into the drain circuit equation (16) so that one must solve for the transmission-line length  $\Theta$  in the drain circuit. The length  $\theta$  must satisfy extended-resonance conditions ( $\theta$  equals either  $\gamma$  or  $\gamma + \pi$ , depending on the sign of  $jB$  at the drain of the

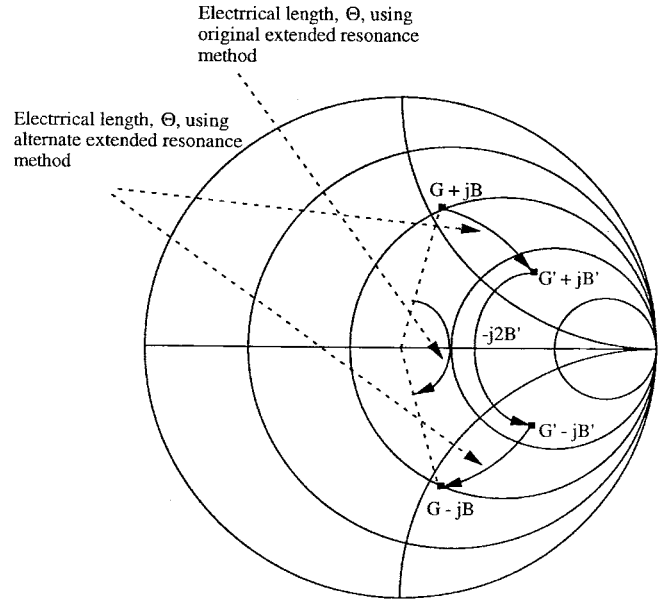


Fig. 7. Smith-chart representation of extended resonance.

device). Assuming a preset value for both  $G$  and  $Y_o$  in (16), one must, therefore, determine what value of  $jB$  (i.e., the value of the shunt reactive element) at the drain of each device must be in order to produce a length  $\theta$  that will satisfy (16).

Transmission lines in the input circuit may sometimes differ substantially in length from those in the output circuit. It may be necessary to use a modified version of the extended-resonance procedure. For example, consider the simple power-dividing circuit in Fig. 6. The difference between this circuit and that in Fig. 4 is that a shunt reactive element is inserted halfway between the devices. The design procedure for this circuit may be seen on the Smith chart in Fig. 7. Instead of converting the admittance,  $Y = G + jB$  of the first

$$V_1 = \frac{I}{Y_o + G + jB} [1 + \Gamma_2 e^{-j2\theta}] + \frac{I e^{-j\phi}}{Y_o + 3G + jB} [1 + \Gamma_1] e^{-j\theta} \quad (13)$$

$$V_2 = \frac{I e^{-j\phi}}{Y_o + 3G + jB} [1 + \Gamma_1 e^{-j2\theta}] + \frac{I}{Y_o + G + jB} [1 + \Gamma_2] e^{-j\theta} \quad (14)$$

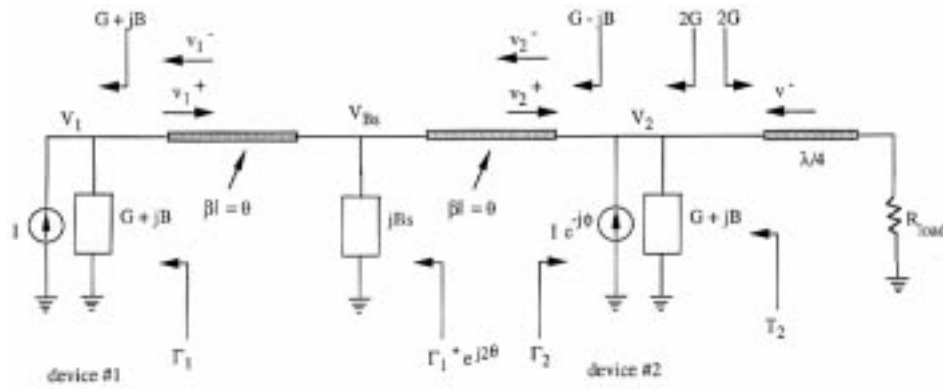


Fig. 8. Analysis of the modified output circuit.

device to its conjugate value  $Y^* = G - jB$ , the admittance of the first device is transformed to some arbitrary value  $Y' = G' + jB'$ , where  $B'$  may be either positive or negative. This transformation is done by the first half of the transmission line. The shunt susceptance  $jB'$ 's convert  $Y' = G' + jB'$  to  $Y'^* = G' - jB'$ . The second half of the transmission line finally takes  $Y'^* = G' - jB'$  to  $Y^* = G - jB$ , to resonate the susceptance of the second device. With a sinusoidal source at the input, the ratio of the voltage at the first device to the voltage at the second device can be determined from (8) and (9). First, the relationship between the voltage at the first device and the voltage halfway between the two devices ( $V_x$ ) is found as follows:

$$\frac{V_1}{V_x} = \frac{[1 + |\Gamma_1|e^{j\gamma}]e^{-j\theta}}{1 + |\Gamma_1|e^{j\gamma}e^{-j2\theta}}. \quad (17)$$

The ratio of  $V_x$  to  $V_2$  is

$$\frac{V_x}{V_2} = \frac{[1 + |\Gamma_1|e^{j\gamma}e^{j2\theta}]e^{-j\theta}}{1 + |\Gamma_1|e^{-j\gamma}}. \quad (18)$$

Thus,  $V_1$  is related to  $V_2$  by

$$\frac{V_1}{V_2} = \frac{[1 + |\Gamma_1|e^{-j\gamma}e^{j2\theta}]e^{-j2\theta}[1 + |\Gamma_1|e^{j\gamma}]}{1 + |\Gamma_1|e^{-j\gamma}[1 + |\Gamma_1|e^{j\gamma}e^{-j2\theta}]}. \quad (19)$$

The magnitude of the voltages  $V_1$  and  $V_2$  are the same as before. In this case,  $V_2$  lags  $V_1$  by

$$\phi = 2\theta - 2 \tan^{-1} \left[ \frac{|\Gamma_1| \sin \gamma}{1 + |\Gamma_1| \cos \gamma} \right] - 2 \tan^{-1} \left[ \frac{|\Gamma_1| \sin(2\theta - \gamma)}{1 + |\Gamma_1| \cos(2\theta - \gamma)} \right]. \quad (20)$$

Fig. 8 illustrates the power-combining scheme using the alternate method. It can be shown that the relationship between  $V_1$  and  $V_2$  for the power-combining version of this circuit is that of (19) with the phase relationship between  $V_2$  and  $V_1$  reversed.

This implies that instead of placing a susceptive stub at the ends of a transmission line inserted between two devices, a single stub may be placed at the center of a shorter length of transmission line that can be inserted between the same two devices. This allows the designer to vary the physical spacing between devices so that it is approximately the same in both the input and output circuits.

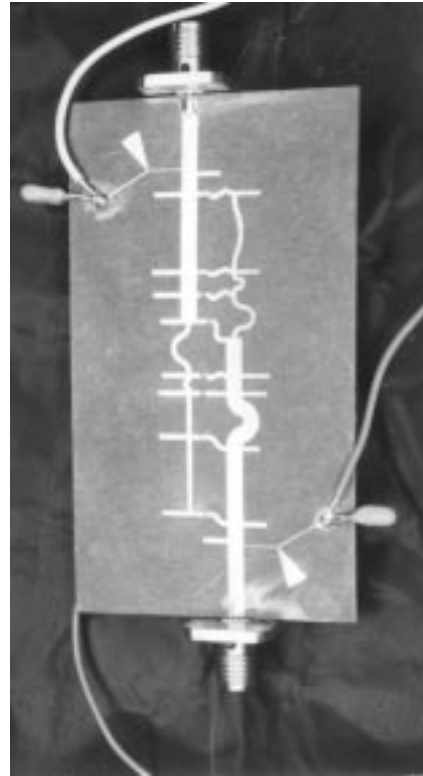


Fig. 9. Photograph of an eight-device extended-resonance amplifier, with the output circuit on top. Circuit measures 73 mm  $\times$  41 mm.

## VI. EXPERIMENT

An eight-device extended-resonance power-combining amplifier was designed and constructed. An enlarged photograph of the eight-device amplifier is shown in Fig. 9. Circuit-board dimensions are 73 mm  $\times$  41 mm. Microstrip stubs provide the desired susceptance at the terminals of each device. The active devices used were 100-mW alpha-power GaAs MESFET chips AF035p1-00 with a gate length of 0.25  $\mu$ m and a total gate periphery of 400  $\mu$ m. Devices were biased at 5 V with a drain current of 70 mA. The amplifier was constructed on a 31-mil-thick Duroid substrate with  $\epsilon_r = 2.33$ . The devices were mounted on copper rods with conductive epoxy. The devices were grounded by inserting the rods through holes in the substrate and soldering them to the ground plane. The gate and drain of each device were each connected to the circuit

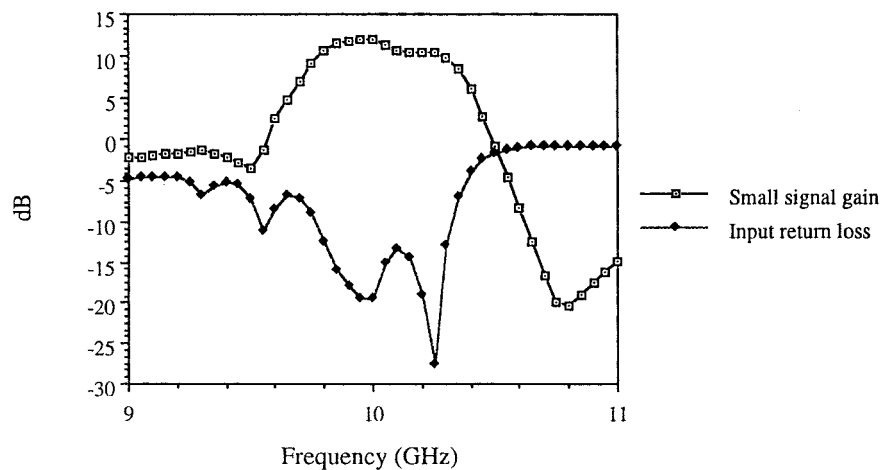


Fig. 10. Predicted small-signal gain and return loss for the eight-device amplifier.

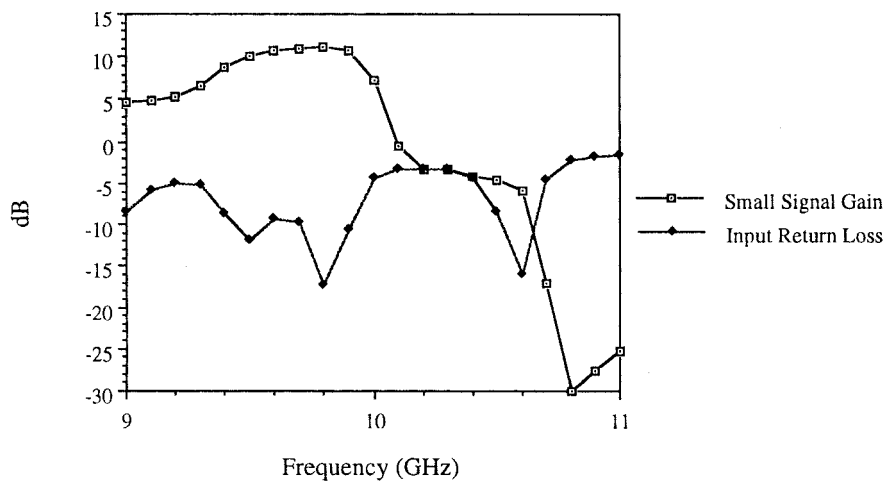


Fig. 11. Measured small-signal gain and return loss for the eight-device amplifier.

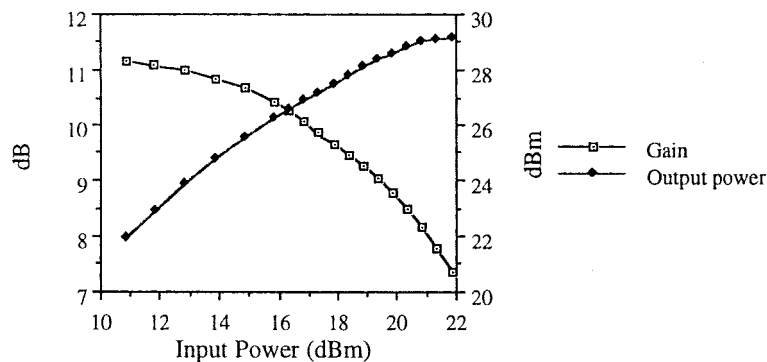


Fig. 12. Measured output power and gain compression curves for the eight-device amplifier.

with bond wire. Fig. 10 shows the predicted small-signal gain and return loss. At 10 GHz, the predicted small-signal gain was 11.9 dB. Fig. 11 shows the measured small-signal gain and return loss. At 9.8 GHz, the measured small-signal gain was 11.2 dB. The predicted and measured performance are very close, except that the measured response is shifted down by about 200 MHz. Otherwise, the measured response is well matched to the predicted performance. The shift in frequency can be due to the inaccuracies in modeling wire bonds and mounting the chips on the substrate. Output power and gain

versus input power at 9.8 GHz is plotted in Fig. 12. At 1-dB compression, the output power is 582 mW with a power-added efficiency of 18.29%. At 3-dB compression, the output power is 813 mW with a power-added efficiency of 30.44%. In order to determine the power-combining performance of the eight-device amplifier, a single-device amplifier was also built and tested. The maximum small-signal gain was 11.5 dB at 9.6 GHz. At the 3-dB compression point, the output power was 111 mW with a power-added efficiency of 33.7%. This indicates a power-combining efficiency of 91.6% for the eight-

device amplifier. The measured 3-dB bandwidth for each amplifier was approximately 5%. As previously mentioned, this bandwidth should meet the demands of many millimeter-wave systems.

## VII. CONCLUSION

An eight-device amplifier based on an extended-resonance technique has been presented. The amplifier was able to produce approximately eight times the power available from a single-device amplifier. Furthermore, a general extended-resonance amplifier design procedure was outlined such that the voltage phase delay in the dividing and combining circuits can be more easily manipulated to permit a combiner circuit with line lengths that are similar to the lengths used in the divider circuit. This general approach alleviates the problem of matching the voltage phase delay through the input and output circuits of the amplifier. Amplifiers based on this more general extended-resonance approach are currently underway. Anticipated goals include the design of a monolithic extended-resonance power-combining amplifier at millimeter-wave frequencies.

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Dr. Mortazawi served as the Orlando IEEE AP/MTT chapter chairman from 1994 to 1995.

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Dr. DeLoach received the IEEE's David Sarnoff Award, and the Franklin Institute's Stuart Ballantine Medal, both in 1975, and the IEEE Medal for Engineering Excellence, in 1993. He has served in many areas of IEEE activities.